

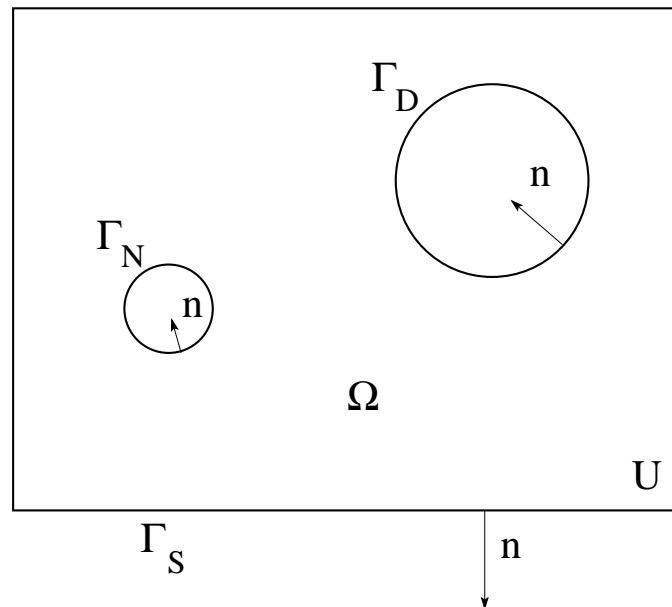
A levelset method in shape optimization for variational inequations

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Introduction

- Ω : Domain in \mathbb{R}^2
- n : Outer normal vector to Ω
- Γ_D : Dirichlet conditions
- Γ_N : Neumann conditions
- Γ_S : Signiorini conditions or Dirichlet conditions



Linear problem

- Linear problem

$$\left\{ \begin{array}{l} -\Delta u + u = f \quad \text{in } \Omega \subset U \subset \mathbb{R}^2 \\ u = 0 \quad \text{on } \Gamma_S \\ u = 0 \quad \text{on } \Gamma_D, \\ \partial_n u = 0 \quad \text{on } \Gamma_N, \end{array} \right.$$

Equivalently, u is solution of the variational equation

$$\int_{\Omega} \nabla u \cdot \nabla v + uv \, d\Omega = \int_{\Omega} f v \, d\Omega \quad \forall v \in H_0^1(\Omega)$$

$$H_0^1(\Omega) = \{v \in H^1(\Omega) \mid v = 0 \text{ on } \Gamma_D \cup \Gamma_S\}$$

Signiorini Problem

- Signiorini Problem

$$\left\{ \begin{array}{l} -\Delta u + u = f \quad \text{in } \Omega \subset U \subset \mathbb{R}^2 \\ u \geq 0, \quad \partial_n u \geq 0, \quad u \partial_n u = 0 \quad \text{on } \Gamma_S \\ u = 0 \quad \text{on } \Gamma_D, \\ \partial_n u = 0 \quad \text{on } \Gamma_N, \end{array} \right.$$

Equivalently, u is the solution of a variational inequality

$$\int_{\Omega} \nabla u \cdot \nabla (v - u) d\Omega \geq \int_{\Omega} (f - u)(v - u) d\Omega \quad \forall v \in K$$

$$K = \{v \in H^1(\Omega) \mid v = 0 \text{ on } \Gamma_D, v \geq 0 \text{ on } \Gamma_S\}$$

- We want to maximize the functional $J(\Omega)$ defined by

$$J(\Omega) = E(\Omega) + \lambda A(\Omega) - \mu P_c(\Omega)^2,$$

with $A(\Omega)$, $P_c(\Omega)$ and $E(\Omega)$ defined by

$$E(\Omega) = \frac{1}{2} \int_{\Omega} |\nabla u|^2 + \frac{1}{2} \int_{\Omega} u^2 - \int_{\Omega} f u$$

$$= -\frac{1}{2} \int_{\Omega} f u,$$

$$A(\Omega) = |\Omega|,$$

$$P_c(\Omega) = \max(0, \mathcal{H}^1(\partial\Omega) - c),$$

Existence of an optimal domain for the linear problem

Theorem 1. (*Bucur and Varchon*)

Let $\Omega_i = U \setminus \bar{\omega}_i$ be a sequence of domains such that the number of connex components of $\bar{\omega}_i$ is uniformly bounded. If $\bar{\omega}_i$ converges to $\bar{\omega}$ in the sense of the Hausdorff metric, then the solution u_i of the Neumann problem

$$\begin{cases} -\Delta u_i + u_i & = f & \text{in } \Omega_i, \\ u_i & = 0 & \text{on } \partial U. \\ \partial u_i / \partial n & = 0 & \text{on } \partial \omega_i. \end{cases}$$

converges to the solution u of the same problem defined on $\Omega = U \setminus \bar{\omega}$ in the following sense :

$$u_i \xrightarrow{L^2(U)} u \text{ and } \nabla u_i \xrightarrow{L^2(U; \mathbb{R}^2)} \nabla u$$

if and only if $|\Omega_i| \rightarrow |\Omega|$. (The functions are implicitly extended by 0 over U)

Theorem 2. *The problem*

$$\max_{\Omega \in \mathcal{O}_k} J(\Omega)$$

admits at least one solution $\Omega \in \mathcal{O}_k = \{\Omega = U \setminus \bar{\omega}, \omega \subset V \subset\subset U, \#\bar{\omega} \leq k\}$.

Proof. Let Ω_i be a maximizing sequence such that $\Omega_i \xrightarrow{H^c} \Omega$. The number of connex components of Ω^c is bounded and

$$\mathcal{H}^1(\partial\Omega) \leq \liminf_i \mathcal{H}^1(\partial\Omega_i).$$

$$\Omega_i \xrightarrow{H^c} \Omega \implies \chi_{\Omega_i} \xrightarrow{L^1(U)} \chi_{\Omega}.$$

As a consequence, we get

$$|\Omega_i| \rightarrow |\Omega|.$$

Finally

$$\begin{aligned} |E(\Omega_i) - E(\Omega)| &= \left| \int_{\Omega_i} f u_i - \int_{\Omega} f u \right| \\ &= \left| \int_U f u_i \chi_{\Omega_i} - \int_U f u \chi_{\Omega} \right| \end{aligned}$$

$$\begin{aligned}
|E(\Omega_i) - E(\Omega)| &= \left| \int_U f u_i (\chi_{\Omega_i} - \chi_{\Omega}) + \int_U f \chi_{\Omega} (u_i - u) \right| \\
&\leq \|u_i\|_2 \|f(\chi_{\Omega_i} - \chi_{\Omega})\|_2 + \|f \chi_{\Omega}\|_2 \|u_i - u\|_2.
\end{aligned}$$

According to theorem (1)

$$\|u_i\|_2 \rightarrow \|u\|_2 \quad \text{and} \quad \|u_i - u\|_2 \rightarrow 0$$

Finally

$$\|f(\chi_{\Omega_i} - \chi_{\Omega})\|_2 \rightarrow 0 \implies E(\Omega_i) \rightarrow E(\Omega),$$

and

$$J(\Omega) \geq \limsup_i J(\Omega_i).$$

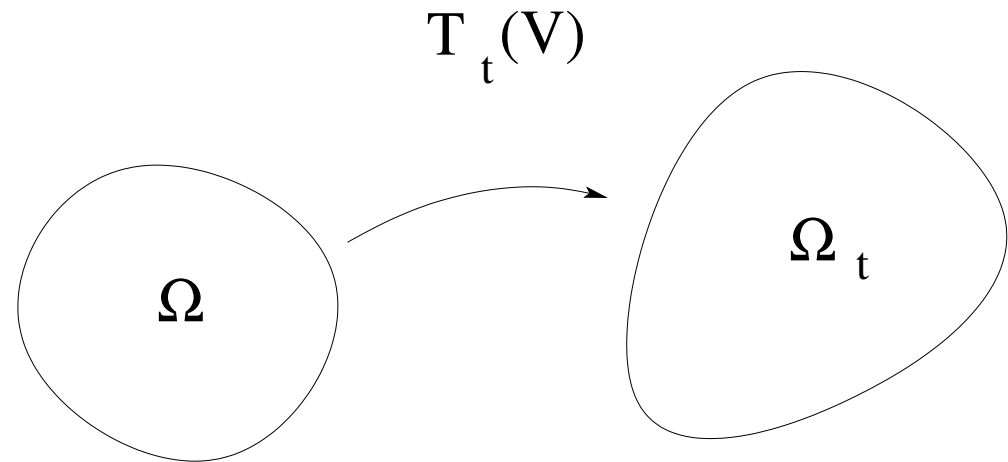
Thus J attains its maximum at Ω .

Shape derivative

Speed method

$V = V(x, t)$ smooth

$$\begin{cases} x'(t) = V(x(t), t), & 0 < t < \tau \\ x(0) = X \in U \subset \mathbb{R}^2 \end{cases}$$



$$T_t(V)(X) = x(t) \text{ and } \Omega_t = T_t(V)(\Omega)$$

- $dJ(\Omega; V) = \left. \frac{dJ(\Omega_t)}{dt} \right|_{t=0} = \lim_{t \rightarrow 0} \frac{J(\Omega_t) - J(\Omega)}{t}$
- $dJ(\Omega; V) > 0 \implies J(\Omega_t) > J(\Omega), \quad t \text{ small.}$

Shape derivative for the linear problem

1) Dirichlet conditions on the boundary of the hole

$$J(\Omega) = E(\Omega) + \lambda A(\Omega) - \mu P_c(\Omega)^2 \quad \text{with} \quad \begin{cases} -\Delta u + u = f & \text{in } \Omega \\ u = 0 & \text{on } \Gamma_S \\ u = 0 & \text{on } \Gamma_D \end{cases}$$

$$dJ(\Omega; V) = \int_{\Gamma_D} \left(-\frac{1}{2} |\nabla u|^2 + \lambda - 2\mu P_c(\Omega) \mathcal{H} \right) V \cdot n \, d\sigma$$

Choice of V_n : we choose $V_n = V \cdot n$ so that $dJ(\Omega; V) > 0$.

$$\longrightarrow \bullet \quad \text{Dirichlet : } V_n = -\frac{1}{2} |\nabla u|^2 + \lambda - 2\mu P_c(\Omega) \mathcal{H} \text{ on } \Gamma_D$$

2) Neumann conditions on the boundary of the hole

$$J(\Omega) = E(\Omega) + \lambda A(\Omega) - \mu P_c(\Omega)^2 \quad \text{with} \quad \begin{cases} -\Delta u + u = f & \text{in } \Omega \\ u = 0 & \text{on } \Gamma_S \\ \partial_n u = 0 & \text{on } \Gamma_N \end{cases}$$

$$dJ(\Omega; V) = \int_{\Gamma_N} \left(\frac{1}{2} |\nabla u|^2 + \frac{1}{2} u^2 - uf + \lambda - 2\mu P_c(\Omega) \mathcal{H} \right) V \cdot n \, d\sigma$$

Choice of V_n : we choose $V_n = V \cdot n$ so that $dJ(\Omega; V) > 0$.

$$\longrightarrow \bullet \quad \text{Neumann : } V_n = \frac{1}{2} |\nabla u|^2 + \frac{1}{2} u^2 - uf + \lambda - 2\mu P_c(\Omega) \mathcal{H} \text{ on } \Gamma_N$$

Shape derivative for the nonlinear problem

Consider transformation $F_\delta = I + \delta V$ with $\Omega_\delta = F_\delta(\Omega)$ and $u^\delta \in K_\delta$ solution of the following variational inequation

$$\int_{\Omega_\delta} \langle \nabla u^\delta, \nabla v - \nabla u^\delta \rangle \geq \int_{\Omega_\delta} (f - u^\delta)(v - u^\delta) \quad \forall v \in K_\delta,$$

where

$$K_\delta = \{v \in H^1(\Omega_\delta) \mid v \geq 0 \text{ on } \Gamma_S ; v = 0 \text{ on } \Gamma_D\} .$$

Energy fonctionnal

$$E(\Omega_\delta) = \frac{1}{2} \int_{\Omega_\delta} |\nabla u^\delta|^2 + (u^\delta)^2 - \int_{\Omega_\delta} f u^\delta.$$

The following limit

$$E'(\Omega, V) := \left. \frac{dE(\Omega_\delta)}{d\delta} \right|_{\delta=0} = \lim_{\delta \rightarrow 0^+} \frac{E(\Omega_\delta) - E(\Omega)}{\delta}$$

exists and is finite, linear and continuous with respect to the vector field V .

$$E(\Omega_\delta) = \frac{1}{2} \int_{\Omega} (\|(DF_\delta^T)^{-1} \nabla u_\delta\| + u_\delta^2) q_\delta dy - \int_{\Omega} f_\delta u_\delta q_\delta dy,$$

with

$$q_\delta = \det DF_\delta = 1 + \delta \operatorname{div}(V) + \delta^2 \det DV .$$

We define

$$f^\delta = q_\delta f_\delta.$$

The following limit exists

$$f' = \left. \frac{df^\delta}{d\delta} \right|_{\delta=0} = \lim_{\delta \rightarrow 0^+} \frac{f^\delta - f^0}{\delta} = \operatorname{div}(V f) \in L^\infty(U) .$$

F_δ bijective between the convex sets K_0 and K_δ and

$$\|u_\delta - u\|_{H^1(\Omega)} \leq C\delta$$

We introduce the notations

$$\pi(\Omega; u) := E(\Omega) = \frac{1}{2} \int_{\Omega} |\nabla u|^2 + u^2 dy - \int_{\Omega} f u dy,$$

$$\pi_\delta(\Omega; u_\delta) := E(\Omega_\delta) = \frac{1}{2} \int_{\Omega} (\|(DF_\delta^T)^{-1} \nabla u_\delta\| + u_\delta^2) q_\delta dy - \int_{\Omega} f_\delta u_\delta q_\delta dy.$$

We get

$$\frac{E(\Omega_\delta) - E(\Omega)}{\delta} = \frac{\pi_\delta(\Omega; u_\delta) - \pi(\Omega; u)}{\delta} \leq \frac{\pi_\delta(\Omega; u) - \pi(\Omega; u)}{\delta},$$

and passing to the limit we find

$$\begin{aligned} \limsup_{\delta \rightarrow 0^+} \frac{E(\Omega_\delta) - E(\Omega)}{\delta} &\leq \lim_{\delta \rightarrow 0^+} \frac{\pi_\delta(\Omega; u) - \pi(\Omega; u)}{\delta} \\ &\leq -\frac{1}{2} \int_{\Omega} \langle A(V) \nabla u, \nabla u \rangle - u^2 \operatorname{div}(V) \, dy - \int_{\Omega} g(V) u \, dy, \end{aligned}$$

with $A(V)$ and $g(V)$ defined by

$$A(V) = DV + DV^T - (\operatorname{div}V)I,$$

$$g(V) = \operatorname{div}(fV).$$

In the same way, we have another inequality

$$\frac{E(\Omega_\delta) - E(\Omega)}{\delta} = \frac{\pi_\delta(\Omega; u_\delta) - \pi(\Omega; u)}{\delta} \geq \frac{\pi_\delta(\Omega; u_\delta) - \pi(\Omega; u_\delta)}{\delta},$$

and passing to the limit, we obtain

$$\begin{aligned} \liminf_{\delta \rightarrow 0^+} \frac{E(\Omega_\delta) - E(\Omega)}{\delta} &\geq \lim_{\delta \rightarrow 0^+} \frac{\pi_\delta(\Omega; u_\delta) - \pi(\Omega; u_\delta)}{\delta} \\ &\geq -\frac{1}{2} \int_{\Omega} \langle A(V) \nabla u, \nabla u \rangle - u^2 \operatorname{div}(V) \, dy - \int_{\Omega} g(V) u \, dy, \end{aligned}$$

thanks to the strong convergence of u_δ to u in $H^1(\Omega)$. Thus

$$\begin{aligned} E'(\Omega, V) &:= \lim_{\delta \rightarrow 0} \frac{E(\Omega_\delta) - E(\Omega)}{\delta} \\ &= -\frac{1}{2} \int_{\Omega} \langle A(V) \nabla u, \nabla u \rangle - u^2 \operatorname{div}(V) \, dy - \int_{\Omega} g(V) u \, dy. \end{aligned}$$

Topological derivative

The speed method does not allow topological changes (creation of a hole).

→ *topological derivative (used by Nazarov, Sokolowski, Zochowski, Guillaume, Masmoudi, Amstutz, Feijoo, Novotny, Tarocco and many others).*

$$\omega_\varepsilon = B_{x_0}(\varepsilon), \quad \Omega_\varepsilon = \Omega \setminus \overline{\omega_\varepsilon}$$

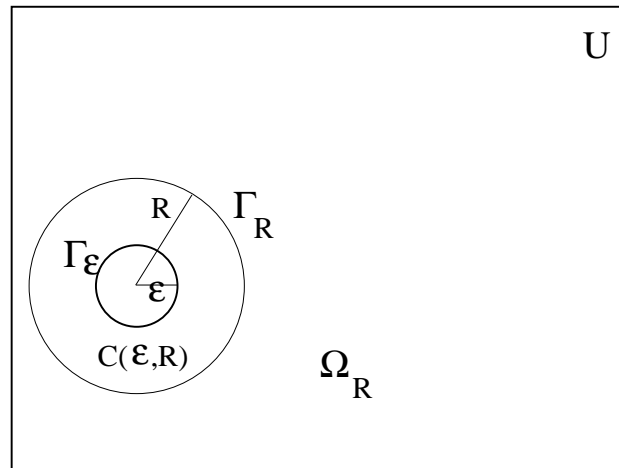
$$J(\Omega_\varepsilon) = -\frac{1}{2} \int_{\Omega_\varepsilon} |\nabla u_\varepsilon|^2 + u_\varepsilon^2 + \lambda A(\Omega_\varepsilon) - \mu P_c(\Omega_\varepsilon)^2$$

Asymptotic expansion :

$$J(\Omega_\varepsilon) = J(\Omega) + \rho(\varepsilon) \mathcal{T}_\Omega(x_0) + o(\rho(\varepsilon)),$$
$$\rho(\varepsilon) > 0, \quad \rho(\varepsilon) \rightarrow 0.$$

If $\mathcal{T}_\Omega(x_0) > 0$ then $J(\Omega_\varepsilon) > J(\Omega)$.

Truncated problem



$$\begin{cases} -\Delta u_\epsilon + u_\epsilon = f & \text{in } \Omega_\epsilon, \quad \Omega_\epsilon = U \setminus B_\epsilon, \\ u_\epsilon = 0 & \text{on } \Gamma_\epsilon. \end{cases}$$

$$\begin{cases} -\Delta u_\epsilon^R + u_\epsilon^R = f & \text{in } \Omega_R, \quad \Omega_R = U \setminus B_R, \quad R > \epsilon, \\ u_\epsilon^R = 0 & \text{on } \Gamma_S, \\ -\partial_n y_\epsilon + \partial_n u_\epsilon^R = A_\epsilon(u_\epsilon^R) & \text{on } \Gamma_R. \end{cases}$$

u_ε^R is solution of the following variational inequality

$$u_\varepsilon^R \in H_{\Gamma_S}^1(\Omega_R) : a_\varepsilon(u_\varepsilon^R, v) = l_\varepsilon(v) \quad \forall v \in H_{\Gamma_S}^1(\Omega_R)$$

$$a_\varepsilon(u, v) = \int_{\Omega_R} \langle \nabla u, \nabla v \rangle + uv \, dS + \int_{\Gamma_R} A_\varepsilon(u)v$$

$$l_\varepsilon(v) = \int_{\Omega_R} f v - \int_{\Gamma_R} v \partial_n y_\varepsilon,$$

Theorem 3. u_ε^R verifies

$$u_\varepsilon^R = u_\varepsilon|_{\Omega_R},$$

The energy of the problem defined on Ω_ε is given by

$$\begin{aligned} E_\varepsilon(f) &= \frac{1}{2} \int_{\Omega_\varepsilon} \|\nabla u_\varepsilon\|^2 + u_\varepsilon^2 \, dS - \int_{\Omega_\varepsilon} f u_\varepsilon \, dS \\ &= -\frac{1}{2} \int_{\Omega_\varepsilon} \|\nabla u_\varepsilon\|^2 + u_\varepsilon^2 \, dS. \end{aligned}$$

Steklov-Poincare operator

The homogeneous problem

We introduce the homogeneous problem in the ring $C(\varepsilon, R) = B(R) \setminus \overline{B(\varepsilon)}$

$$\begin{cases} -\Delta w_\varepsilon + w_\varepsilon = 0 & \text{in } C(\varepsilon, R), \\ w_\varepsilon = v & \text{on } \Gamma_R, \\ w_\varepsilon = 0 & \text{on } \Gamma_\varepsilon. \end{cases}$$

The Steklov-Poincare operator is defined by

$$A_\varepsilon : H^{1/2}(\Gamma_R) \rightarrow H^{-1/2}(\Gamma_R)$$

$$A_\varepsilon(v) = \frac{\partial w_\varepsilon}{\partial n} \quad \text{on } \Gamma_R$$

Asymptotic expansion for A_ε

$$A_\varepsilon = A_0 + \frac{1}{|\ln \varepsilon|} B + o\left(\frac{1}{|\ln \varepsilon|^2}\right) \text{ in } \mathcal{L}(H^{1/2}(\Gamma_R); H^{-1/2}(\Gamma_R))$$

We define the energies

$$E^{(1)}(v) = \int_{B(R)} \|\nabla w\|^2 + w^2 dS, \quad E_\varepsilon^{(1)}(v) = \int_{C(R,\varepsilon)} \|\nabla w_\varepsilon\|^2 + w_\varepsilon^2 dS,$$

Theorem 4. $E_\varepsilon^{(1)}(v)$ has an expansion

$$E_\varepsilon^{(1)}(v) = E^{(1)}(v) + \frac{\pi a_0(R)^2}{2I_0(R)^2 |\ln \varepsilon|} + \mathfrak{R}(v),$$

with

$$|\mathfrak{R}(v)| \leq \frac{M}{|\ln \varepsilon|^2},$$

uniformly on bounded subsets of $H^{\frac{1}{2}}(\Gamma_R)$ and

$$v(R, \phi) = \frac{1}{2}a_0(R) + \sum_{k=1}^{\infty} a_k(R) \sin k\phi + b_k(R) \cos k\phi$$

$$I_k(r) = \sum_{m=0}^{\infty} \frac{\left(\frac{r}{2}\right)^{k+2m}}{m!(k+m)!}, \quad k \geq 0$$

The non-homogeneous problem

$$\begin{cases} -\Delta y_\varepsilon + y_\varepsilon = f|_{C(R,\varepsilon)} & \text{in } C(R,\varepsilon) \\ y_\varepsilon = 0 & \text{on } \Gamma_R \\ y_\varepsilon = 0 & \text{on } \Gamma_\varepsilon \end{cases}$$

with $f \in C^\infty(\mathbb{R}^2)$. We study the following function :

$$f|_{C(R,\varepsilon)} \mapsto \frac{\partial y_\varepsilon}{\partial n}|_{\Gamma_R} = g_\varepsilon.$$

Theorem 5. g_ε admits the expansion

$$g_\varepsilon = g_0 - \frac{h(R)}{2RI_0(R) \ln \varepsilon} + O((\ln \varepsilon)^{-2}),$$

with

$$h(R) = \frac{1}{\pi} \int_0^R \int_0^{2\pi} t f(t, \phi) \frac{I_0(t)K_0(R) - K_0(t)I_0(R)}{I_0(R)} dt d\phi.$$

We define the energies

$$E^{(2)}(f) = - \int_{B(R)} \|\nabla y\|^2 + y^2 dS, \quad E_\varepsilon^{(2)}(f) = - \int_{C(R,\varepsilon)} \|\nabla y_\varepsilon\|^2 + y_\varepsilon^2 dS$$

Theorem 6. $E_\varepsilon^{(2)}(f)$ has an expansion

$$E_\varepsilon^{(2)}(f) = E^{(2)}(f) + \frac{\pi h(R)^2}{2|\ln \varepsilon|} + \mathfrak{R}(f),$$

with

$$|\mathfrak{R}(f)| \leq \frac{M}{|\ln \varepsilon|^2},$$

Theorem 7. *The energy of the problem defined on Ω_ε admits the expansion*

$$E_\varepsilon(f) = E(f) + \frac{\pi a_0(R)^2}{4I_0(R)^2 |\ln \varepsilon|} + \frac{\pi h(R)^2}{4 |\ln \varepsilon|} + \frac{\pi a_0 h(R)}{2I_0(R) |\ln \varepsilon|} + \mathfrak{R}(f)$$

with

$$|\mathfrak{R}(f)| \leq \frac{M}{|\ln \varepsilon|^2}$$

where the coefficient $a_0(R)$ is defined by

$$a_0(R) = \frac{1}{\pi} \int_0^{2\pi} u(R, \phi) d\phi.$$

We have the following convergences

$$\lim_{R \rightarrow 0} h(R) = 0, \quad \lim_{R \rightarrow 0} a_0(R) = 2u(0), \quad \lim_{R \rightarrow 0} I_0(R) = 1$$

so that we can consider the following approximation for the topological derivative

$$\mathcal{T}_\Omega(0) = \frac{1}{4\pi} \left(\int_0^{2\pi} u(R, \phi) d\phi \right)^2$$

In fact, we can prove that

$$u(0) = \frac{a_0(R)}{2I_0(R)} + \frac{h(R)}{2}$$

As a consequence

$$E_\varepsilon(f) = E(f) + \frac{2\pi u(0)^2}{|\ln \varepsilon|} + \mathfrak{R}(f).$$

with

$$\mathfrak{R}(f) \leq \frac{M}{|\ln \varepsilon|^2}$$

- Dirichlet conditions on $\partial\omega_\varepsilon$

$$u_\varepsilon = 0 \text{ on } \partial\omega_\varepsilon$$

$$J(\Omega_\varepsilon) = J(\Omega) + \frac{\pi u(0)^2}{|\ln \varepsilon|} + o(1/\log \varepsilon) - 4\mu P_c(\Omega)\pi\varepsilon - \lambda\pi\varepsilon^2$$

$$\rightarrow \mathcal{T}_\Omega(x_0) = u_\Omega^2(x_0) \geq 0$$

- Neumann conditions on $\partial\omega_\varepsilon$

$$\partial_n u_\varepsilon = 0 \text{ on } \partial\omega_\varepsilon$$

$$J(\Omega_\varepsilon) = J(\Omega) + \pi\varepsilon^2 \left(-|\nabla u(x_0)|^2 - \frac{1}{2}u(x_0)^2 + uf(x_0) - \lambda \right) + 4\mu P_c(\Omega)\pi\varepsilon + o(\varepsilon^2)$$

$$\rightarrow \begin{cases} \mathcal{T}_\Omega(x_0) = -|\nabla u(x_0)|^2 - \frac{1}{2}u(x_0)^2 + uf(x_0) - \lambda & \text{si } P_c(\Omega) = 0 \\ \mathcal{T}_\Omega(x_0) = 4\mu P_c(\Omega)\pi & \text{si } P_c(\Omega) > 0 \end{cases}$$

If $\mathcal{T}_\Omega(x_{max}) \geq \mathcal{T}_\Omega(x)$ for all $x \in \Omega$ in $\mathcal{T}_\Omega(x_{max}) > 0$, we drill a hole in Ω at x_{max} .

Signiorini problem in Ω_R

$$\left\{ \begin{array}{ll} -\Delta u_\varepsilon^R + u_\varepsilon^R = f & \text{in } \Omega_R, \quad \Omega_R = U \setminus B_R, \quad R > \varepsilon, \\ \partial_n u_\varepsilon^R = 0 & \text{on } \Gamma_N, \\ u_\varepsilon^R \geq 0, \quad \partial_n u_\varepsilon^R \geq 0, \quad u_\varepsilon^R \partial_n u_\varepsilon^R = 0 & \text{on } \Gamma_S \\ -\partial_n y_\varepsilon + \partial_n u_\varepsilon^R = A_\varepsilon(u_\varepsilon^R) & \text{on } \Gamma_R, \end{array} \right.$$

verifies $u_\varepsilon^R = u_R + \varepsilon^2 q_R + o(\varepsilon^2)$ in $H^1(\Omega_R)$. We also have

$$\begin{aligned} u_\varepsilon^R &= u_\varepsilon|_{\Omega_R} \quad \forall \varepsilon \geq 0 \\ q_R &= q|_{\Omega_R} \end{aligned}$$

The outer topological derivative q of the solution u is given by

$$\int_{\Omega} \nabla q \nabla (v - q) + q(v - q) + \int_{\Gamma_R} B(u)(v - q) \geq 0$$

where B is the derivative of the Steklov-Poincare operator and

$$q \in S(u) = \{v \in H^1(\Omega) \mid v = 0 \text{ on } \Gamma_D, v \geq 0 \text{ on } \Xi(u), a(u, v) = 0\}$$

$$a(u, v) = \int_{\Omega} \nabla u \nabla v + uv - \int_{\Omega} f v$$

$$\Xi(u) = \{x \in \Gamma_s \mid u(x) = 0\}$$

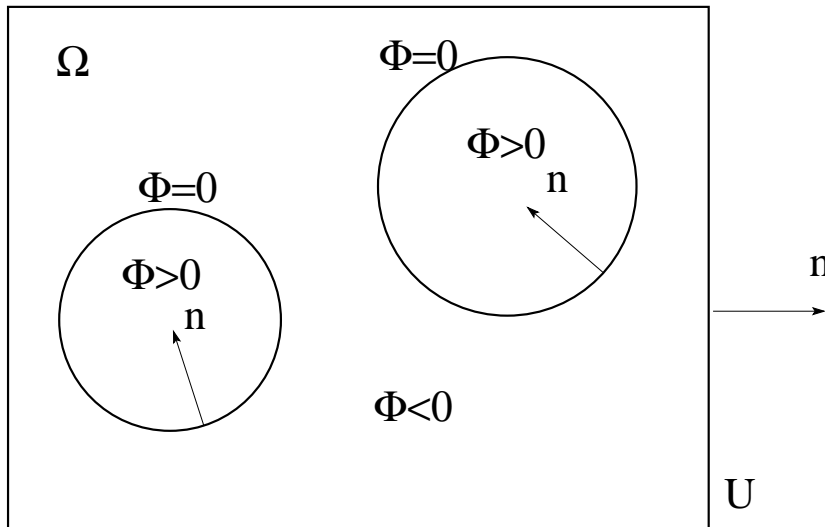
We also obtain an expression similar to the linear case for the topological derivative of the functional J .

The levelset function

The levelset function $\phi = \phi(x, t)$ is defined for all $x \in U$

$$\Omega = \{x \in U, \phi(x, t) < 0\}$$

$$\Gamma = \{x \in U, \phi(x, t) = 0\}$$



$$\left. \begin{aligned} \phi_t + V \cdot \nabla \phi &= 0 \text{ in } U \\ n &= \frac{\nabla \phi}{|\nabla \phi|} \end{aligned} \right\}$$

$$\Rightarrow \text{Hamilton-Jacobi equation} \\ \phi_t + V_n |\nabla \phi| = 0 \text{ in } U$$

- Dirichlet : $V_n = -\frac{1}{2}|\nabla u|^2 + \lambda - 2\mu P_c(\Omega)\mathcal{H}$ on Γ_D
- Neumann : $V_n = \frac{1}{2}|\nabla u|^2 + \frac{1}{2}u^2 - uf + \lambda - 2\mu P_c(\Omega)\mathcal{H}$ on Γ_N

Remarks : V_n must be defined everywhere in U , not only on $\Gamma \rightarrow$ extension
We want to have $|\nabla \phi| \simeq 1$.

Extension of the normal speed

We calculate V_{ext} in U so that $V_{ext} = V_n$ on $\Gamma = \Gamma_D$ or Γ_N and

$$\nabla V_{ext} \cdot \nabla \phi = 0 \text{ in } U.$$

Then we get $|\nabla \phi| = 1$ for all t .

Construction de V_{ext}

Solve (for large T)

$$\begin{aligned} q_\tau + S(\phi) \frac{\nabla \phi}{|\nabla \phi|} \cdot \nabla q &= 0, & \text{in } U \times (0, T) \\ q(x, \tau = 0) &= p(x), & x \in U \end{aligned}$$

where $p = V_n$ on Γ and $p = 0$ otherwise.

S is close to the sign function

$$S(d) = \frac{d}{\sqrt{d^2 + |\nabla d|^2 \epsilon^2}}, \quad \epsilon = \min(\Delta x, \Delta y).$$

Algorithm

- Initialisation of ϕ^0 , initial domain Ω_0 .
- For k from 0 to K
 - (1) Calculation of the topological derivative in Ω_k
 - i. calculation of u_{Ω_k} in Ω_k
 - ii. calculation of the topological derivative
 - iii. a hole is created \rightarrow domain Ω_k^1
 - (2) Solve Hamilton-Jacobi equation in U
 - * For p from 1 to P
 - i. calculation of $u_{\Omega_k^p}$ in Ω_k^p
 - ii. calculation of V_n^p on Γ_k^p and its extension to U
 - iii. solve Hamilton-Jacobi equation : $\phi_t + V_n^p |\nabla \phi| = 0$, \rightarrow domain Ω_k^{p+1}
 - end for p
 - * $\rightarrow \phi^{k+1}$ in D and new domain $\Omega_{k+1} = \Omega_k^P$
- end for k

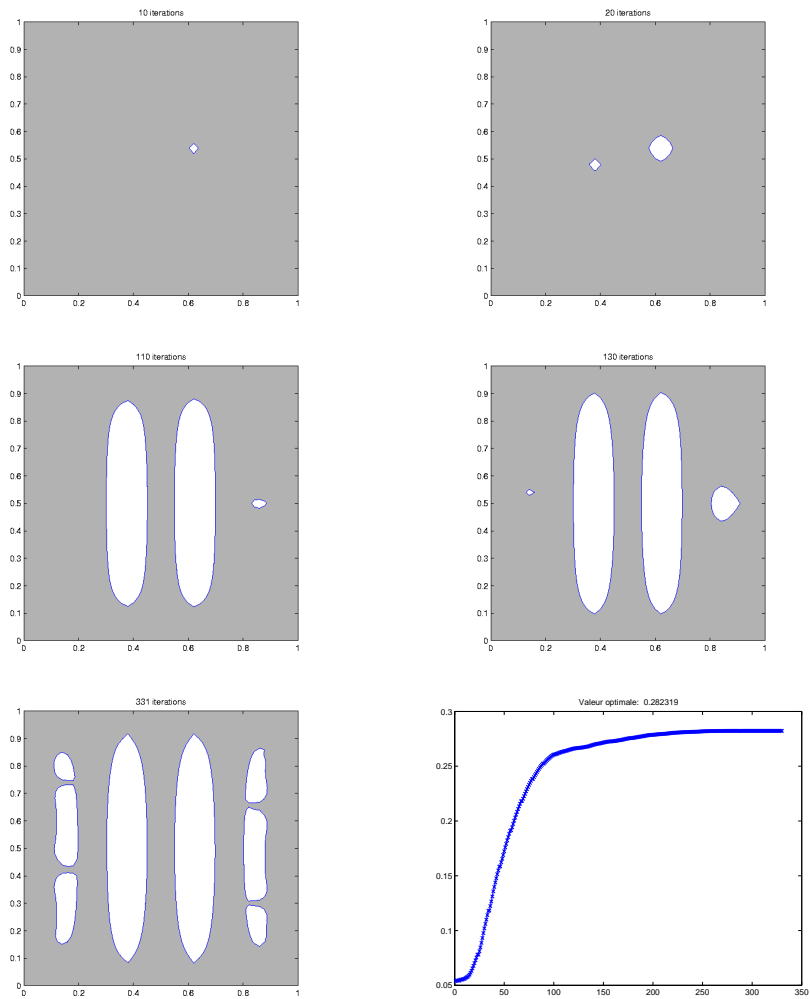


FIG. 1. The initial domain is full, we use the topological derivative

$$f = 10 \sin^2(4\pi x), \lambda = 0.5, \mu = 0$$

Optimal value : $J(\Omega_1) = 0,282319$

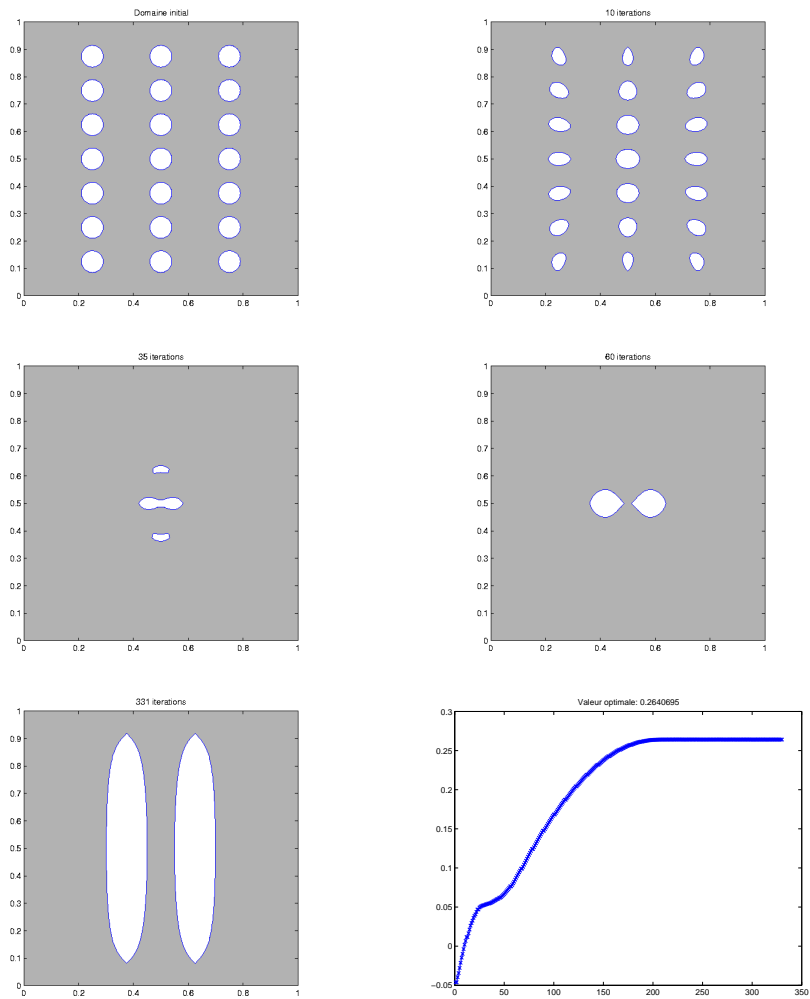


FIG. 1. 21 holes in the initial domain, no topological derivative

$$f = 10 \sin^2(4\pi x), \lambda = 0.5, \mu = 0$$

Optimal value : $J(\Omega_2) = 0,2640695$

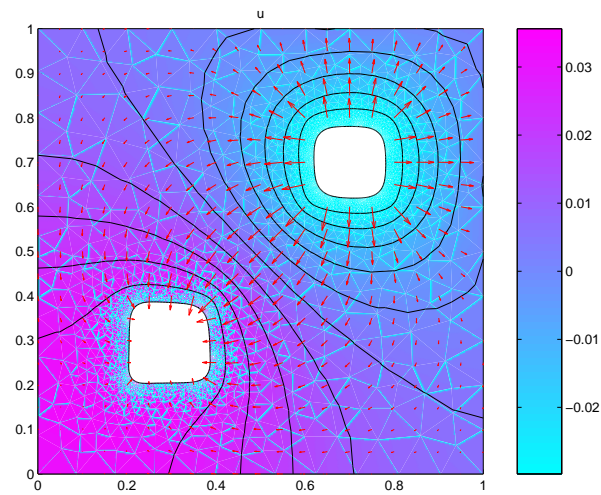
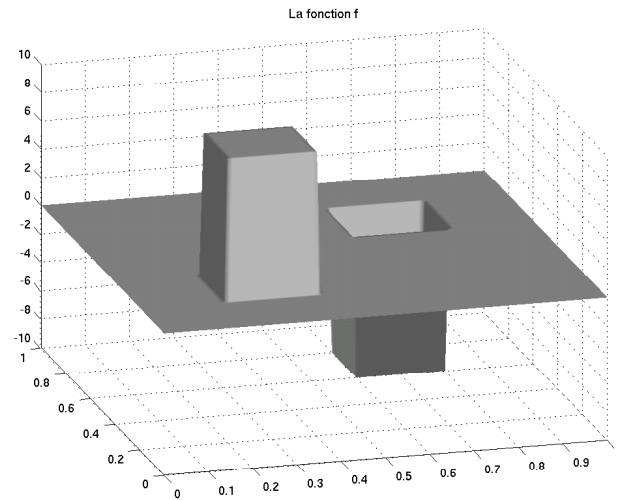


FIG. 1. Function f and solution u in the final domain

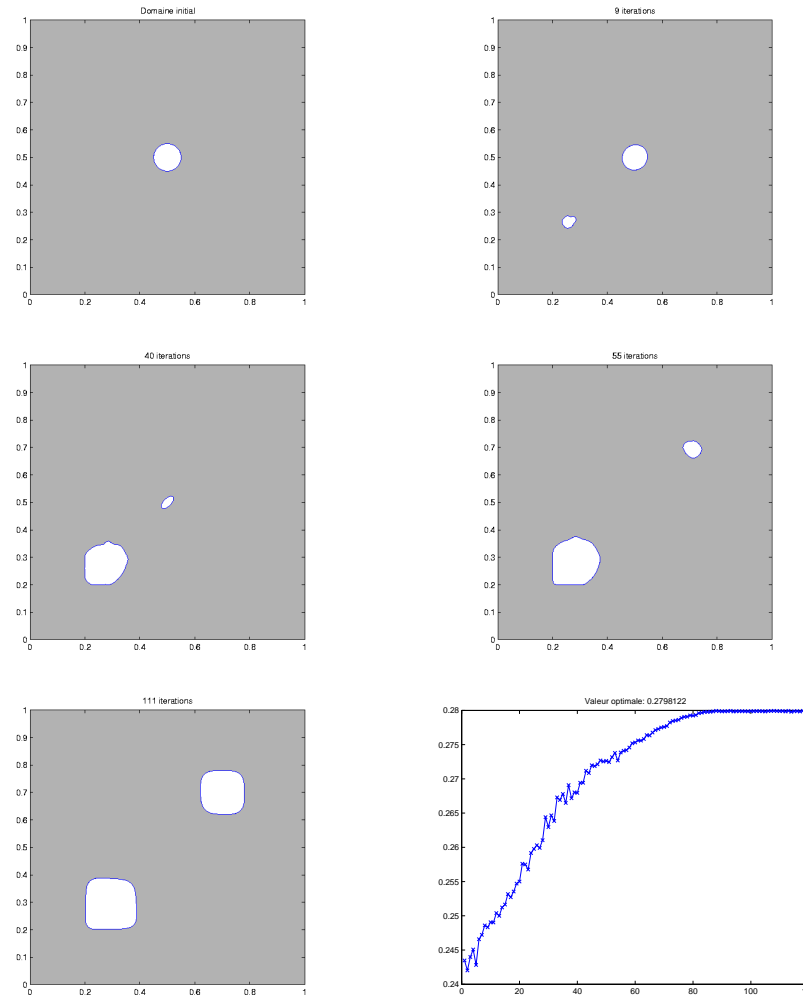


FIG. 1. The non-linear case
 $\lambda = 0.3, \mu = 0.001, c = 0.6$

Conclusion

- Without topological derivative, the algorithm is very sensitive to the initial domain, a good initial guess is necessary.
- With topological derivative, we can start with any initial domain, in particular the full domain.
- Heavy calculations without topological derivative, due to an initial domain with a lot of holes.
- For the non-linear problem, it is possible to define a shape derivative and a topological derivative, and we obtain the same formulas as for the linear problem.